

Aminodimethylalane (Me₂AlNH₂): Matrix Isolation and *ab Initio* Calculations

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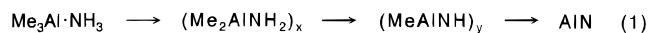
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Abstract: The calculated harmonic frequencies (RMP2(fc)/6-31G*) of the three isomeric alanes Me₃Al·NH₃, Me₃Al·¹⁵NH₃, and Me₃Al·ND₃ are assigned to the measured IR absorptions of these compounds, matrix isolated in argon at 10 K. The results are compared with the known literature data. The irradiations of the matrix-isolated amine adducts yield aminodimethylalane. The experimental IR data of the three isotopomers Me₂AlNH₂, Me₂Al¹⁵NH₂, and Me₂AlND₂ agree well with the calculated harmonic frequencies. Elucidated from *ab initio* calculations, Me₂AlNH₂ is a C_{2v}-symmetrical molecule with an AlN bond length of 1.790 Å. This AlN distance is 15% shorter than that of the starting compound Me₃Al·NH₃ (2.106 Å). The predicted AlN rotation barrier of 40.7 kJ mol⁻¹ is much lower than the known barriers of the isoelectronic silaalkenes and the lighter congeners, aminoboranes and alkenes, respectively.

Introduction

The first aminoalane was prepared by Wiberg over half a century ago.¹ He investigated the pyrolysis of Me₃Al·NH₃, which yielded the amino- and iminoalane and, finally, aluminium nitride. Either the alkane elimination or the simple “salt



synthesis”, starting from alkali metal amides and aluminium halides, are the common preparation methods for aminoalanes. To date, a large variety of compounds with the general formula R₂Al(NR'R'') are known. They usually exist as dimers or trimers.² Power et al. isolated the first monomeric derivatives using sterically demanding substituents.³ The primary focus of Power's work had been the study of a π-interaction between aluminium and nitrogen. Based on ¹H-NMR experiments, it was possible to estimate an AlN rotation barrier of 38–42 kJ mol⁻¹ for a certain derivative.³ Other examples of unassociated alanes with known structural details are Al[N(SiMe₃)₂]₃, Al-(NiPr₂)₃, and MesAl[N(SiMe₃)₂]₂.⁴ These substances exhibit trigonal-planar coordinated Al and N atoms. However, the two respective planes are twisted in the AlN direction and there is no striking correlation between the torsion angles and the AlN bond lengths. Indeed, the bulky ligands provide monomeric species, necessary to examine a possible π-interaction, but the steric effects may be of the same order of magnitude as the electronic effects one is interested in and both effects might be indistinguishable.

This report describes the synthesis and characterization of

[⊗] Abstract published in *Advance ACS Abstracts*, June 15, 1996.

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aminodimethylalane (Me₂AlNH₂) in an argon matrix, the first monomeric aminoalane without sterically demanding substituents.

Results and Discussion

Starting Compounds. The three isomeric alanes Me₃Al·NH₃, Me₃Al·¹⁵NH₃, and Me₃Al·ND₃ are readily accessible by the reaction of trimethylalane with ammonia at low temperature. Subsequent sublimation (30 °C, 10⁻³ mbar) gives the pure compounds in high yields (≥95%). Though ammonia trimethylalane has been known for more than five decades¹ and has been reinvestigated several times, the purification and characterization by standard methods has not been described in the literature.⁵ We checked the purity of the alanes by NMR (¹H, ¹³C, ²⁷Al), MS, and CHN elemental analysis. As expected, there is a broad resonance in the ¹³C-NMR spectra and a singlet in the ²⁷Al-NMR spectra for each isotopomer. The ²⁷Al-NMR shift of δ 168 is in the typical range for tetracoordinated aluminium with these kinds of ligands.⁶ The isotopic labeling is revealed by the ¹H-NMR measurements: two singlets for Me₃Al·NH₃, one singlet for the methyl groups with a sharp doublet for the H₃¹⁵N protons of Me₃Al·¹⁵NH₃, and only one singlet for Me₃Al·ND₃. Corresponding results are obtained from the mass spectra where M⁺ – CH₃ with *m/e* 74, 75, and 77 for Me₃Al·NH₃, Me₃Al·¹⁵NH₃, and Me₃Al·ND₃, respectively, is the heaviest fragment.

With respect to our *ab initio* calculation at the RMP2(fc)/6-31G* level, ammonia trimethylalane has C_{3v} symmetry (Figure 1).⁷ The AlC₃ skeleton of the molecule deviates only slightly from planarity (C–Al–C = 117.3°). The AlN distance of 2.106 Å is relatively long compared to the common range of simple

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(6) Akitt, J. W. In *Annual Reports on NMR Spectroscopy*; Mooney, E. F., Ed.; Academic Press: London, New York, 1972; Vol. 5A, p 545.

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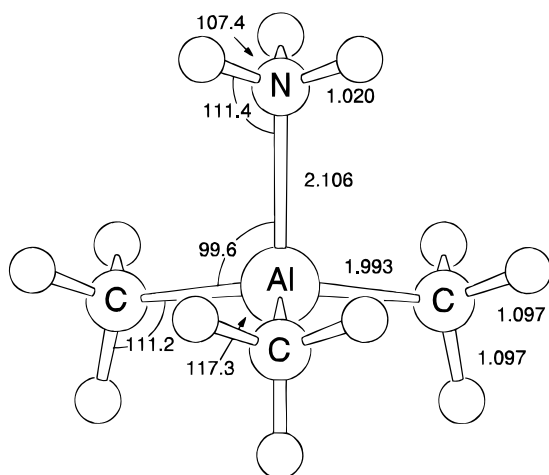


Figure 1. Selected bond distances (Å) and angles (deg) for the equilibrium geometry of Me₃Al·NH₃ (*C*_{3v} point group) calculated at the RMP2(fc)/6-31G* level.

amine adducts, which are known to have AlN bond lengths between 1.94 and 2.10 Å.⁸ The experimental structure of Me₃Al·NH₃ is unknown, however, the isomeric molecule H₃Al·NMe₃ is well-known in the literature.⁹ Elucidated from gas-phase electron diffraction, this aluminium hydride exhibits *C*_{3v} symmetry with a flat AlH₃ pyramid (H–Al–H = 114.1°) and an AlN bond length of 2.063 Å.^{9a} These results are in accord with high-level *ab initio* calculations published recently (e.g., AlN = 2.063 Å, H–Al–H = 117.1° at the CISD/DZP level).^{9c} The authors searched for other possible equilibrium geometries by using different methyl-group conformations, but only one configuration gave energy minima on the potential surfaces at several levels of theory. As expected, our calculation for Me₃Al·NH₃ reveals a similar *C*_{3v}-symmetrical equilibrium geometry, especially concerning the methyl-group conformations (Figure 1).

Using the conventional matrix-isolation technique (Experimental Section) we prepared argon matrices of the three isotopomers Me₃Al·NH₃, Me₃Al·¹⁵NH₃, and Me₃Al·ND₃. Figure 2b–d depicts the corresponding IR spectra measured at 10 K in the range of 250–4000 cm⁻¹.

Ammonia trimethylalane is expected to have 10 A₁, 5 A₂, and 15 E modes, in which the A₂ modes are IR inactive. For comparison with experiment we calculated the harmonic vibrational frequencies of the adducts at the RMP2(fc)/6-31G* level.⁷ Table 1 compiles these data and Figure 2a depicts the scaled theoretical IR spectrum of Me₃Al·NH₃. Based on these predictions we assigned the detected IR bands and the results are given in Table 2.

Ault previously investigated matrix-isolated ammonia trimethylalane.¹⁰ He used a merged jet of Me₃Al with NH₃ to obtain Me₃Al·NH₃ in argon matrices. He then assigned seven IR absorptions of both Me₃Al·NH₃ and Me₃Al·¹⁵NH₃ and eight of Me₃Al·ND₃ (Table 2, numbers in braces). Though we confirmed most of his findings, we disagree on two points. In accord with our calculation, the CH₃-rocking modes ν₂₂ are detected at 731.5, 731.0, and 720.0 cm⁻¹ for Me₃Al·NH₃, Me₃Al·¹⁵NH₃, and Me₃Al·ND₃, respectively, whereas Ault assigned the same value (733 cm⁻¹) to each of the three isotopomers.

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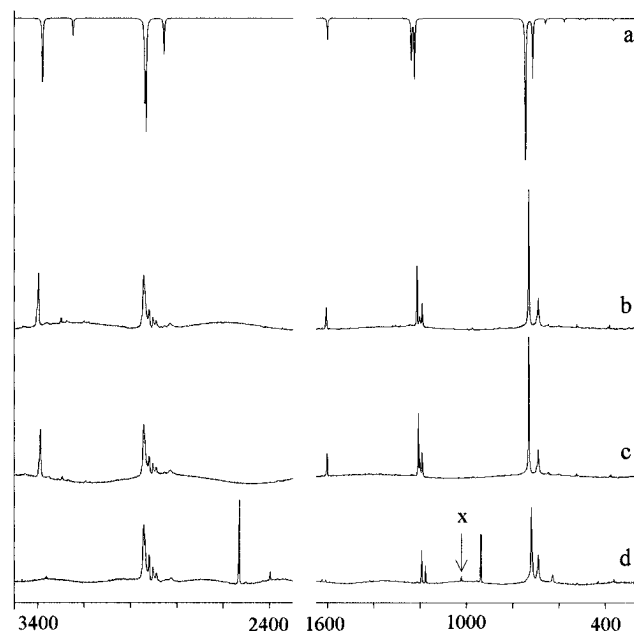


Figure 2. IR spectra of ammonia trimethylalane: X-axis, 250–1650 cm⁻¹ and 2300–3500 cm⁻¹; Y-axis, Absorbance (dimensionless). The absorptions of the 2300–3600-cm⁻¹ region are enhanced by a factor of 3 relative to the 250–1650-cm⁻¹ range. Spectrum a: Me₃Al·NH₃ calculated at the RMP2(fc)/6-31G* level (frequencies scaled by 0.93, half band width of 2 cm⁻¹). Spectrum b: Me₃Al·NH₃ in argon at 10 K. Spectrum c: Me₃Al·¹⁵NH₃ in argon at 10 K. Spectrum d: Me₃Al·ND₃ in argon at 10 K. x (1022.0 cm⁻¹) indicates small amounts of Me₃Al·ND₂H.¹⁰

Table 1. Harmonic Frequencies (cm⁻¹) and Intensities (in Parentheses, km mol⁻¹) Calculated at the RMP2(fc)/6-31G* Level

	Me ₃ Al·NH ₃	Me ₃ Al· ¹⁵ NH ₃	Me ₃ Al·ND ₃
A ₁ : ν ₁	178.1 (11.8)	177.2 (11.5)	175.3 (11.4)
ν ₂	392.8 (6.8)	386.0 (7.2)	371.3 (5.6)
ν ₃	521.8 (4.4)	521.8 (4.4)	521.7 (4.2)
ν ₄	768.8 (180.2)	768.8 (180.2)	768.8 (181.0)
ν ₅	1317.8 (102.3)	1314.3 (160.2)	1322.4 (2.7)
ν ₆	1331.9 (129.1)	1328.0 (68.1)	1014.8 (136.9)
ν ₇	1531.2 (1.4)	1531.2 (1.4)	1531.2 (1.4)
ν ₈	3071.1 (0.2)	3071.1 (0.2)	3071.1 (0.2)
ν ₉	3160.0 (74.7)	3160.0 (74.7)	3160.0 (74.7)
ν ₁₀	3490.7 (18.3)	3488.3 (17.6)	2496.1 (15.1)
A ₂ : ν ₁₁	82.5 (0.0)	82.5 (0.0)	59.3 (0.0)
ν ₁₂	113.8 (0.0)	113.8 (0.0)	112.9 (0.0)
ν ₁₃	596.3 (0.0)	596.3 (0.0)	596.3 (0.0)
ν ₁₄	1521.8 (0.0)	1521.8 (0.0)	1521.8 (0.0)
ν ₁₅	3151.4 (0.0)	3151.4 (0.0)	3151.4 (0.0)
E: ^a ν ₁₆	117.3 (1.7)	117.2 (1.8)	115.6 (3.4)
ν ₁₇	131.6 (9.8)	130.5 (9.0)	124.8 (4.6)
ν ₁₈	171.3 (0.4)	171.2 (0.5)	168.3 (1.3)
ν ₁₉	550.8 (3.9)	549.5 (4.4)	454.4 (8.4)
ν ₂₀	621.7 (10.5)	621.5 (10.0)	616.9 (5.8)
ν ₂₁	708.7 (17.2)	707.2 (19.3)	674.8 (41.3)
ν ₂₂	802.0 (432.4)	801.3 (431.4)	788.8 (345.4)
ν ₂₃	1317.3 (86.5)	1317.3 (86.5)	1317.3 (83.3)
ν ₂₄	1518.4 (1.9)	1518.4 (1.9)	1518.5 (1.5)
ν ₂₅	1529.9 (0.9)	1529.9 (0.9)	1530.3 (2.2)
ν ₂₆	1721.1 (68.1)	1718.1 (67.4)	1246.2 (39.6)
ν ₂₇	3070.1 (36.1)	3070.1 (36.1)	3070.1 (36.1)
ν ₂₈	3152.9 (108.7)	3152.9 (108.7)	3152.9 (109.1)
ν ₂₉	3158.7 (7.0)	3158.7 (7.0)	3158.7 (7.1)
ν ₃₀	3633.2 (65.2)	3623.3 (63.7)	2675.1 (43.9)

^a Intensities of the E modes are the sum of the two degenerated vibrations.

Secondly, we cannot confirm the assignment of the IR band at 627 cm⁻¹ of the deuterated compound to the asymmetric AlC₃-stretching mode ν₂₀. This band is more likely the degenerate

Table 2. Experimental Frequencies (cm⁻¹) and Observed/Calculated Frequency Ratios (in Parentheses)^a

		Me ₃ Al·NH ₃	Me ₃ Al· ¹⁵ NH ₃	Me ₃ Al·ND ₃
A ₁ : ν ₂	ν(AIN)	383.0 (0.9751)	376.5 (0.9754)	364.0 (0.9803)
E: ν ₁₉	δ'(NH ₃) _{rock}	517.5 (0.9395)	516.5 (0.9399)	432.0 (0.9507)
A ₁ : ν ₃	ν _s (AlC ₃)	523.0 (1.0023)	523.0 (1.0023)	523.0 (1.0025)
		{523}	{523}	{523}
E: ν ₂₀	ν _{as} (AlC ₃)	601.0 (0.9667)	600.5 (0.9662)	
E: ν ₂₁	δ'(CH ₃) _{rock}	647.5 (0.9136)	646.5 (0.9142)	627.0 (0.9292)
A ₁ : ν ₄	δ'(CH ₃) _{rock}	690.0 (0.8975)	690.0 (0.8975)	689.5 (0.8969)
E: ν ₂₂	δ'(CH ₃) _{rock}	731.5 (0.9121)	731.0 (0.9123)	720.0 (0.9128)
		{733}	{733}	{733}
A ₁ : ν ₅	δ _s (CH ₃)	1191.0 (0.9038)	1191.0 (0.9062)	
E: ν ₂₃	δ _s (CH ₃)	1203.0 (0.9132)	1201.5 (0.9121)	1192.5 (0.9053)
A ₁ : ν ₆	δ _s (NH ₃)	1213.0 (0.9107)	1207.5 (0.9093)	937.0 (0.9233)
		{1212}	{1207}	{936}
E: ν ₂₆	δ _{as} (NH ₃)	1605.5 (0.9328)	1602.0 (0.9324)	1176.5 (0.9473)
		{1606}	{1602}	{1175}
A ₁ : ν ₈ - ν ₉	ν(CH ₃)	2943.0, 2936.0	2942.5, 2936.0	2943.0, 2936.0
E: ν ₂₇ - ν ₂₉		2920.0, 2904.0	2920.0, 2903.0	2919.5, 2903.5
		2888.0	2888.5	2887.5
		{2942; 2919}	{2942; 2919}	{2942; 2919}
A ₁ : ν ₁₀	ν _s (NH ₃)	3296.5 (0.9444)	3291.5 (0.9436)	2397.0 (0.9603)
		[3300.0 (0.9454)]	[3295.5 (0.9447)]	[2399.0 (0.9611)]
E: ν ₃₀	ν _{as} (NH ₃)	3395.5 (0.9346)	3387.5 (0.9349)	2530.0 (0.9458)
		[3402.5 (0.9365)]	[3402.5 (0.9391)]	[2534.0 (0.9473)]
		{3396}	{3388}	{2530}

^a Wavenumbers in braces were taken from the literature.¹⁰ Absorptions due to matrix effects are shown in square brackets.

CH₃-rocking mode ν₂₁; the calculated isotopic shifts and the intensities were nicely reproduced by our experiments.¹¹ Besides, the asymmetric stretch ν₂₀ is predicted to be of low intensity; in particular, this vibration should be much weaker for Me₃Al·ND₃ compared with the other two isotopomers. We tentatively assign the very weak absorptions at 601.0 and 600.5 cm⁻¹ for Me₃Al·NH₃ and Me₃Al·¹⁵NH₃, respectively, to ν₂₀; we could not detect a signal for the deuterated molecule.

The most interesting low-energy vibration is the AlN-stretching mode ν₂ at 383.0, 376.5, and 364.0 cm⁻¹ (Table 2). An early vibrational study on ammonia trimethylalane assigned the weak IR bands at 461 and 439 cm⁻¹, measured at -196 °C for the solid samples of Me₃Al·NH₃ and Me₃Al·ND₃, respectively, to the AlN-stretching mode.^{5a,12} This assignment had been confirmed by Raman spectroscopy and normal coordinate analysis. The differences between these and our values may be due to different media: the pure solids there and highly diluted matrices here. The isotopic shifts caused by the deuteration are the same in the solid state and in the matrix (solid: 439/461 = 0.9523; matrix: 364.0/383.0 = 0.9504), suggesting that we are really looking at the same modes in both cases. A comment should be made on the CH₃ stretches. The calculated values do not reproduce the measured ones exactly (Figure 2). Nevertheless, as predicted by the *ab initio* calculation, no differences between the CH₃ stretches of all three isotopomers should be seen, as confirmed by the experiments.

In summary, we found all expected normal modes between 250 and 4000 cm⁻¹ except the three weak vibrations ν₇, ν₂₄, and ν₂₅ (Table 1) and the asymmetric AlC₃-stretching mode ν₂₀ of Me₃Al·ND₃ already discussed.

Photolysis of Ammonia Trimethylalane. The irradiation of matrix-isolated ammonia trimethylalane with UV light at 210 nm¹³ yields methane and the hitherto unknown monomeric aminodimethylalane.



Figure 3 shows a typical IR spectrum obtained after 7 h of

(11) The weak vibrations between 250 and 650 cm⁻¹ cannot be seen easily in Figure 2, but they are readily revealed in a higher magnification.



Figure 3. Irradiation of Me₃Al·NH₃ in argon at 10 K: X-axis, 250–1650 and 2800–3600 cm⁻¹; Y-axis, Absorbance (dimensionless). Spectrum a: Me₂AlNH₂ calculated at the RMP2(fc)/6-31G* level (frequencies scaled by 0.93, half band width of 2 cm⁻¹). Spectrum b: Difference IR spectrum after irradiation (210 nm)¹³ of Me₃Al·NH₃ for 7 h (decreasing absorptions point down, increasing absorptions point up). x indicates CH₄ (1304.0, 1306.5, 3024.0 (br) cm⁻¹), y indicates NH₃ (974.5 cm⁻¹), and z indicates H₂O (1624.0 cm⁻¹).

irradiation of Me₃Al·NH₃ at 10 K (Figure 3b), with the scaled theoretical spectrum of Me₂AlNH₂ (Figure 3a).⁷ Under the

(12) The AlN-stretching mode was out of Ault's¹⁰ detectable range of 400–4000 cm⁻¹.

(13) UV light with the maximum at 210 nm and a band width of 20 nm was used (see Experimental Section).

(14) DCH₃ has another resonance at 1156.0 cm⁻¹ which is hidden by a strong absorption of Me₂AlND₂ (ν₇, Table 4). Herzberg, G. *Molecular Spectra and Molecular Structure*, 1st ed.; D. Van Nostrand Company, Inc.: Princeton, NJ, 1945; Vol. II, p 309.

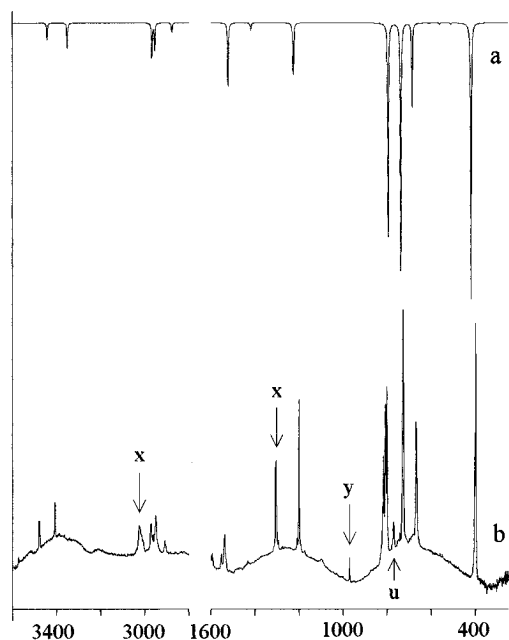


Figure 4. IR spectra of Me₂Al¹⁵NH₂: X-axis, 250–1600 and 2800–3600 cm⁻¹; Y-axis, Absorbance (dimensionless). Spectrum a: Me₂Al¹⁵NH₂ calculated at the RMP2(fc)/6-31G* level (frequencies scaled by 0.93, half band width of 2 cm⁻¹). Spectrum b: Me₂Al¹⁵NH₂ in argon at 10 K synthesized by the irradiation (20 h, 210 nm¹⁵) of Me₃Al·¹⁵NH₃ (Figure 2c). x indicates CH₄ (1304.5, 3024.0 (br) cm⁻¹), y indicates ¹⁵NH₃ (970.5 cm⁻¹), and u indicates an unassigned product band (770.0 cm⁻¹).

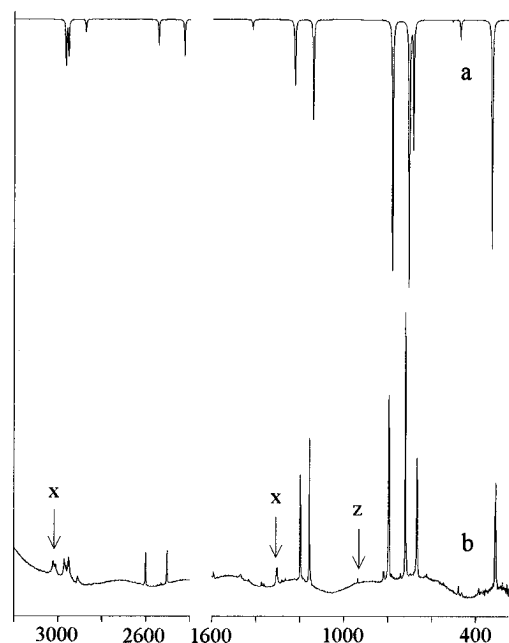


Figure 5. IR spectra of Me₂AlND₂: X-axis, 250–1600 and 2400–3200 cm⁻¹; Y-axis, Absorbance (dimensionless). Spectrum a: Me₂AlND₂ calculated at the RMP2(fc)/6-31G* level (frequencies scaled by 0.93, half band width of 2 cm⁻¹). Spectrum b: Me₂AlND₂ in argon at 10 K synthesized by the irradiation (20 h, 210 nm¹⁵) of Me₃Al·¹⁵ND₃ (Figure 2d). x indicates DCH₃ (1303.5, 3010.5, 3020.5 cm⁻¹)¹⁴ and z indicates small amounts of Me₃Al·ND₃ (937.0 cm⁻¹).

conditions employed for our experiments, the methane loss is complete after 20 h; only traces of the starting material are detectable. The corresponding IR spectra of the ¹⁵N-labeled sample and of the deuterated compound with their scaled calculated spectra are depicted in Figures 4 and 5, respectively.

Ab initio calculations predict a C_{2v}-symmetry equilibrium geometry for aminodimethylalane (Figure 6).⁷ Consequently,

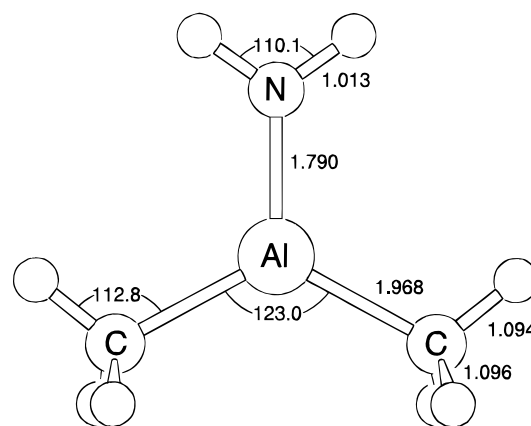


Figure 6. Selected bond distances (Å) and angles (deg) for the equilibrium geometry of Me₂AlNH₂ (C_{2v} point group) calculated at the RMP2(fc)/6-31G* level.

Table 3. Harmonic Frequencies (cm⁻¹) and Intensities (in Parentheses, km mol⁻¹) Calculated at the RMP2(fc)/6-31G* Level

	Me ₂ AlNH ₂	Me ₂ Al ¹⁵ NH ₂	Me ₂ AlND ₂
A ₁ : ν ₁	164.0 (2.7)	163.5 (2.6)	163.1 (2.6)
ν ₂	551.0 (1.4)	548.6 (1.1)	545.6 (1.0)
ν ₃	743.5 (8.1)	738.6 (4.6)	727.8 (1.4)
ν ₄	863.0 (181.0)	855.0 (183.8)	839.4 (165.1)
ν ₅	1319.1 (17.4)	1319.1 (17.3)	1319.2 (15.0)
ν ₆	1523.8 (1.3)	1523.8 (1.3)	1524.1 (2.3)
ν ₇	1641.2 (57.4)	1635.7 (54.7)	1226.9 (63.6)
ν ₈	3093.1 (1.9)	3093.1 (1.9)	3093.1 (1.9)
ν ₉	3189.9 (19.4)	3190.0 (19.4)	3189.9 (19.5)
ν ₁₀	3608.8 (24.0)	3603.6 (22.8)	2610.4 (23.9)
A ₂ : ν ₁₁	56.5 (0.0)	56.5 (0.0)	56.4 (0.0)
ν ₁₂	402.3 (0.0)	402.3 (0.0)	291.7 (0.0)
ν ₁₃	655.7 (0.0)	655.7 (0.0)	644.4 (0.0)
ν ₁₄	1515.3 (0.0)	1515.3 (0.0)	1515.3 (0.0)
ν ₁₅	3176.3 (0.0)	3176.3 (0.0)	3176.3 (0.0)
B ₁ : ν ₁₆	187.5 (2.5)	185.6 (2.5)	170.3 (1.8)
ν ₁₇	604.9 (1.1)	603.7 (1.3)	506.1 (12.9)
ν ₁₈	654.0 (0.7)	653.0 (0.5)	641.3 (0.0)
ν ₁₉	795.9 (214.0)	793.4 (212.3)	758.7 (175.2)
ν ₂₀	1316.5 (37.4)	1316.5 (37.4)	1316.5 (37.1)
ν ₂₁	1519.9 (0.0)	1519.9 (0.0)	1519.8 (0.1)
ν ₂₂	3092.2 (6.2)	3092.2 (6.2)	3092.2 (6.2)
ν ₂₃	3190.0 (10.8)	3190.0 (10.8)	3190.0 (10.8)
ν ₂₄	3712.3 (16.0)	3701.8 (15.0)	2736.5 (16.2)
B ₂ : ν ₂₅	68.2 (0.1)	68.1 (0.1)	67.8 (0.1)
ν ₂₆	195.3 (6.8)	194.7 (7.0)	192.3 (4.4)
ν ₂₇	453.6 (239.5)	450.7 (236.2)	351.8 (146.5)
ν ₂₈	738.1 (71.4)	738.1 (70.9)	737.4 (81.9)
ν ₂₉	1525.2 (5.1)	1525.2 (5.1)	1525.2 (4.9)
ν ₃₀	3176.9 (23.7)	3176.9 (23.7)	3176.9 (23.7)

one expects 30 normal modes belonging to the irreducible representations A₁ (10), A₂ (5), B₁ (9), and B₂ (6). Table 3 gives a summary of the calculated harmonic frequencies for the three isomers Me₂AlNH₂, Me₂Al¹⁵NH₂, and Me₂AlND₂. The A₂ modes are IR inactive, and the wavenumbers for ν₁, ν₁₆, ν₂₅, and ν₂₆ are too low to be detectable (<250 cm⁻¹). Therefore, 21 fundamentals remain in our accessible range (250–4000 cm⁻¹). Based on the *ab initio* calculation, we assigned 11 resonances in the IR spectra of Me₂AlNH₂ and Me₂Al¹⁵NH₂, and 12 IR bands for Me₂AlND₂ to 14 and 15 normal modes, respectively (Table 4). The remaining unassigned fundamentals, which were calculated to be very weak (Table 3), have not been measured.

After the photolysis at 10 K the IR bands of aminodimethylalane show fine structure, which is very sensitive to temperature changes. At 20 K we observed a strong relative intensity change for all product vibrations (Figure 7). A second short warmup to 20 K does not yield this annealing effect again. The change

Table 4. Experimental Frequencies (cm^{-1}) and Observed/Calculated Frequency Ratios (in Parentheses)

		Me_2AlNH_2	$\text{Me}_2\text{Al}^{15}\text{NH}_2$	Me_2AlND_2
$\text{B}_2: \nu_{27}$	$\kappa(\text{NH}_2)_{\text{wagg}}$	399.0 (0.8796) [407.0 (0.8973)]	396.5 (0.8797) [404.5 (0.8975)]	307.5 (0.8741) [313.5 (0.8911)]
$\text{B}_1: \nu_{17}$	$\delta'(\text{NH}_2)_{\text{rock}}$			478.0 (0.9448)
$\text{B}_2: \nu_{28}$	$\delta_{\text{as}}(\text{CH}_3)$	668.0 (0.9050)	668.0 (0.9050)	666.5 (0.9039)
$\text{B}_1: \nu_{19}$	$\delta'(\text{CH}_3)_{\text{rock}}$	728.0 (0.9147) [732.0 (0.9197)]	726.5 (0.9157) [730.5 (0.9207)]	720.0 (0.9490)
$\text{A}_1: \nu_4$	$\nu(\text{AlN})^b$	807.5	800.5	794.0 (0.9459)
		817.0	809.0	[796.0 (0.9483)]
		823.5	816.5	
$\text{B}_1: \nu_{20}$	$\delta_{\text{s}}(\text{CH}_3)$	1199.5 (0.9111)	1199.5 (0.9111)	1197.5 (0.9096)
$\text{A}_1: \nu_5$	$\delta_{\text{s}}(\text{CH}_3)$	1199.5 (0.9093)	1199.5 (0.9093)	1197.5 (0.9077)
$\text{A}_1: \nu_7$	$\delta_{\text{s}}(\text{NH}_2)$	1544.0 (0.9408)	1536.0 (0.9390)	1156.0 (0.9422)
$\text{B}_1: \nu_{22}$	$\nu_{\text{s}}(\text{CH}_3)$	2909.0 (0.9408)	2909.0 (0.9408)	2910.0 (0.9411)
$\text{A}_1: \nu_8$	$\nu_{\text{s}}(\text{CH}_3)$	2909.0 (0.9405)	2909.0 (0.9405)	2910.0 (0.9408)
$\text{B}_2: \nu_{30}$	$\nu_{\text{as}}(\text{CH}_3)$	2951.5 (0.9291)	2951.0 (0.9289)	2951.5 (0.9291)
$\text{A}_1: \nu_9$	$\nu_{\text{as}}(\text{CH}_3)$	2971.0 (0.9314)	2971.5 (0.9315)	2970.5 (0.9312)
$\text{B}_1: \nu_{23}$	$\nu_{\text{as}}(\text{CH}_3)$	2971.0 (0.9313)	2971.5 (0.9315)	2970.5 (0.9312)
$\text{A}_1: \nu_{10}$	$\nu_{\text{s}}(\text{NH}_2)$	3411.0 (0.9452)	3406.5 (0.9453)	2503.0 (0.9589)
$\text{B}_1: \nu_{24}$	$\nu_{\text{as}}(\text{NH}_2)$	3486.5 (0.9392)	3478.5 (0.9397)	2600.5 (0.9503)

^a Absorptions due to matrix effects are shown in square brackets. ^b See discussion on the $\nu(\text{AlN})$ of Me_2AlNH_2 and $\text{Me}_2\text{Al}^{15}\text{NH}_2$.

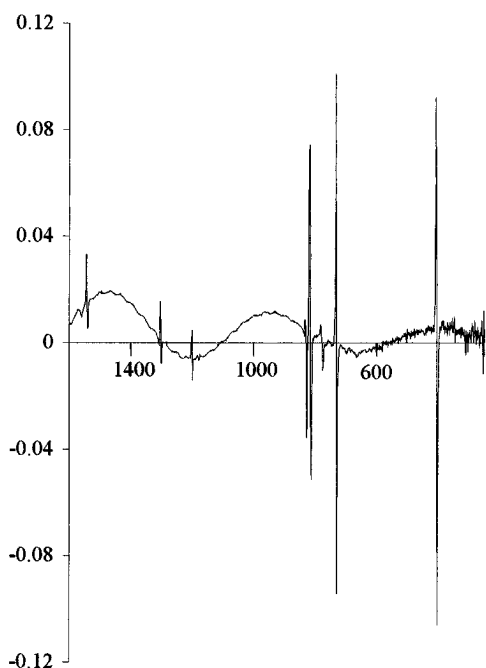


Figure 7. Difference IR spectrum ($250\text{--}1600\text{ cm}^{-1}$) of matrix-isolated Me_2AlNH_2 synthesized by the irradiation (20 h, 210 nm^{13}) of $\text{Me}_3\text{Al}\cdot\text{NH}_3$ (Figure 2a), before (down) and after (up) a warmup from 10 to 20 K.

in intensity initially observed may arise from a release of the hydrostatic pressure built up in the matrix by the reaction (eq 2). Consistent with this conclusion, the methane vibrations are also affected by the warmup, the pressure being exerted on both molecules generated in the argon matrix. Similar hydrostatic pressure effects are known in the literature.¹⁵ Due to matrix effects some bands remain split (Table 4: ν_{27} , ν_{19} , and ν_4 , numbers in brackets), with the largest splitting for the NH_2 -wagging mode ν_{27} ($6.0\text{--}8.0\text{ cm}^{-1}$).

A comment should be made on the assignment of the AlN-stretching vibration ν_4 . In the case of the deuterated compound (Figure 5) the intense absorption at 794.0 cm^{-1} unequivocally corresponds to this stretching mode. Considering the calculated isotopic shifts, the AlN stretches of Me_2AlNH_2 and $\text{Me}_2\text{Al}^{15}\text{NH}_2$ are expected at 816.3 and 808.8 cm^{-1} , respectively. Instead

of one band, three bands were measured in the range of $800\text{--}825\text{ cm}^{-1}$ (Table 4, Figures 3 and 4). This is probably the result of Fermi resonance between ν_4 and the overtone of the intense mode ν_{27} . The situation is complex, because the ν_{27} fundamental, as mentioned earlier, consists of two bands. Simple doubling of the wavenumber of ν_{27} would cause the two overtone couples $798.0/814.0\text{ cm}^{-1}$ (Me_2AlNH_2) and $793.0/809.0\text{ cm}^{-1}$ ($\text{Me}_2\text{Al}^{15}\text{NH}_2$) as a zero approximation. The value of the higher frequency in each couple is very close to calculated value of the AlN-stretching mode ν_4 (816.3 and 808.8 cm^{-1}). In the deuterated molecule the estimated frequencies of the ν_{27} overtones ($615.0/627.0\text{ cm}^{-1}$) are too far away from the ν_4 vibrational level (794.0 cm^{-1}) to allow Fermi resonance.

Besides the frequencies, Table 4 exhibits the observed/calculated frequency ratios (numbers in parentheses). The worst fit has been found for the low-energy NH_2 -wagging mode ν_{27} ($0.8741\text{--}0.8797$) and the best fit for the asymmetric NH_2 stretch ($0.9452\text{--}0.9589$). Within a certain margin of error, one would expect the same ratio for all three isotopomers for a certain mode, if the isotopic shifts had been correctly calculated. Actually, only the fundamental ν_{19} deviates from that expectation for unknown reasons (Table 4).

In summary, we found all normal modes with predicted intensities larger than 10 km mol^{-1} . Nevertheless, two medium strong IR bands at 770.0 and 772.0 cm^{-1} for Me_2AlNH_2 and $\text{Me}_2\text{Al}^{15}\text{NH}_2$, respectively, have not yet been assigned (Figure 4, u). These two absorptions are sensitive to the warmup procedure (Figure 7) discussed earlier. Therefore, we assume them to be absorptions of aminodimethylalane, probably combination vibrations. No comparable vibration was found for the deuterated species.

Reaction 2 represents a photochemical elimination of methane. As expected, this is a 1,2-elimination, which is proven for $\text{Me}_3\text{Al}\cdot\text{ND}_3$ by the formation of DCH_3 (Figure 5). So far, the nature of the electron transitions involved is not known. The corresponding thermolytic version of eq 2 had been investigated first by Wiberg in 1939.¹ Heating of $\text{Me}_3\text{Al}\cdot\text{NH}_3$ in benzene under reflux yielded the trimer of aminodimethylalane, $(\text{Me}_2\text{AlNH}_2)_3$.¹⁶ Sauls et al. examined the kinetics of this process revealing a complex mechanism.^{5b} They proposed that monomeric Me_2AlNH_2 acts as a catalyst. The same authors elucidated the equilibrium constant for the dimer-trimer interconversion from $^1\text{H-NMR}$ experiments ($30\text{--}100\text{ }^\circ\text{C}$) with-

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out detecting the monomer.¹⁷ Recently, Power et al. isolated the first monomeric aminodialkylalanes using bulky ligands.³ With one exception, these compounds show trigonal-planar coordinated Al and N atoms, but the two corresponding planes are twisted along the AlN bond axis. The torsion angles vary between 5.5 and 84.9°, without exhibiting a strong correlation with the AlN bond lengths (1.784–1.880 Å). The *ab initio* calculation for Me₂AlNH₂ reveals a C_{2v} symmetry (Figure 6) with an AlN distance of 1.790 Å. A similar molecular structure, based on valence bond theory, was predicted by Messmer et al. (AlN = 1.797 Å).¹⁸ In a simple MO description, this equilibrium geometry allows a maximum of pπ–pπ interaction between aluminum and nitrogen; hence, it fulfills the expectations for a classical double bond.

It is common to judge single- or double-bond character either by bond lengths or rotation barriers. With a calculated AlN distance of 2.106 Å (Me₃Al·NH₃, Figure 1), the AlN bond shortening in Me₂AlNH₂ becomes 15.0%. This is a large value, if compared with the 9.4% found for the isoelectronic silaalkene Me₂SiCH₂, whose corresponding SiC bond length of 1.692 Å¹⁹ contrasts with the SiC single bond length of 1.868 Å (both values from microwave spectra).²⁰ For the lighter congeners, aminoboranes and alkenes, the corresponding values are 10.8 and 13.6%, respectively.²¹ One has to be careful not to overestimate the 15% shortening in the present case. First, we are comparing calculated bond lengths that may be slightly different from the real ones. Second, the AlN distances of simple amine adducts are known to vary in quite a broad range of 1.94–2.10 Å,^{8,9} making a well-defined distance for a single bond unavailable. Third, it is unknown how much of the shortening goes back to the π-bond interaction, on the one hand, or to the decrease in coordination number from four to three, which usually causes a decrease in the covalent radii, on the other hand.

The optimized transition-state geometry for the AlN rotation (Figure 8) corresponds to a barrier of 40.7 kJ mol⁻¹.²² This is a small value in comparison with silaalkenes (148–149 kJ mol⁻¹ in H₂SiCH₂),²³ aminoboranes (123–139 kJ mol⁻¹ in H₂BNH₂),²⁴ and ethylenes (272 kJ mol⁻¹). The transition state reveals a weak pyramidalization of the NH₂ group and a slight increase of the AlN distance (Figure 8). If the nitrogen lone pair is not involved in any π-bond, the value of 1.823 Å corresponds to a “pure” AlN single bond. This would mean that the π-overlap alone causes a decrease of the AlN distance by 0.033 Å (1.8%). Two recently published *ab initio* calculations of the parent aminoalane H₂AlNH₂ led to similar findings: The AlN bond differences between the ground and the transition states were

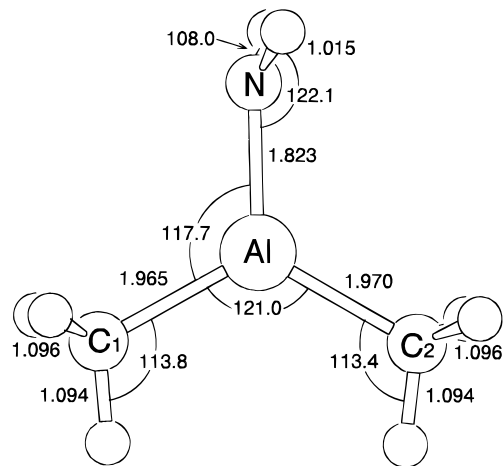


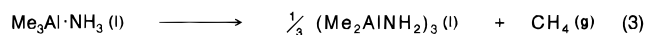
Figure 8. Selected bond distances (Å) and angles (deg) of the transition state (C_s point group) for the AlN rotation of Me₂AlNH₂ calculated at the RMP2(fc)/6-31G* level.

Table 5. Calculated Total Energies *E* (au) and Zero-Point Energies ZPE (Scaled by 0.93 in kJ mol⁻¹) with the Number of Imaginary Frequencies Shown in Parentheses

		<i>E</i>	ZPE
Me ₃ Al·NH ₃	C _{3v}	-417.60357	365.1 (0)
Me ₂ AlNH ₂	C _{2v}	-377.27170	243.3 (0)
Me ₂ AlNH ₂	C _s	-377.25459	239.1 (1)
Me ₃ Al	C _{3h}	-361.20348	267.1 (0)
CH ₄	T _d	-40.33255	113.2 (0)
NH ₃	C _{3v}	-56.35421	86.2 (0)

0.023 (1.3%)²⁵ and 0.019 Å (1.1%)²⁶ with barriers of 46.8 and 44.8 kJ mol⁻¹, respectively. To our knowledge, the only experimental AlN rotation barrier of a monoaminoalane was measured for *t*Bu₂AlNH(Diip) (Diip = 2,6-diisopropylphenyl). Based on ¹H-NMR experiments it was estimated to be 38–42 kJ mol⁻¹,³ in good agreement with our calculation. Certainly, every method for separating a π-bond from a σ-bond includes approximations. In the case of a rotation barrier the pπ–pπ interaction is minimized for the transition state, whereas the same geometry meets the optimum conditions for other π-type overlaps like “negative hyperconjugation”,²⁵ which is —although present—quite probably of minor importance for Me₂AlNH₂. Despite these uncertainties, aminodimethylalane contains a weak AlN double bond, concluded from the rotation barrier.

The *ab initio* calculation for Me₂AlNH₂ yields a heat of reaction (eq 2) of -10.4 kJ mol⁻¹ (Table 5). The experimental value for this methane loss is not known, but the overall decomposition reaction (eq 3) of the neat liquid amine adduct was found to proceed with Δ*H* = -82.2 kJ mol⁻¹.^{5b} The heat of the cyclotrimerization of Me₂AlNH₂ is then approximately -71.8 kJ mol⁻¹.



To check the reliability of the predicted thermodynamic data, we calculated the dissociation enthalpy of the amine adduct. A value of Δ*H* = 93 ± 5 kJ mol⁻¹ was measured for this process

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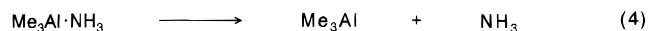
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in benzene (eq 4).^{5b} The heat of complexation of $\text{AlMe}_3(\text{g})$ to



benzene and the heat of solution of $\text{NH}_3(\text{g})$ in benzene are -8 and -12 kJ mol^{-1} , respectively.^{5b} The heat of solution of $\text{Me}_3\text{Al}\cdot\text{NH}_3$ is not known, but it is probably slightly exothermic. Considering these figures, we obtain $\Delta H \leq 113 \pm 5 \text{ kJ mol}^{-1}$ for reaction 4 in the gas phase, which agrees very well with our calculation of $108.6 \text{ kJ mol}^{-1}$ (Table 5).²²

Experimental Section

The matrix apparatus consists of a vacuum line (Leybold Turbovac 151; Leybold Trivac D4B) and a Displex CSW 202 cryogenic closed-cycle system (APD Cryogenics Inc.) fitted with CsI windows. In a typical experiment, argon (Linde 5.6) was deposited for 1–2 h at 15 K (CsI window) with an argon flow of ca. 4 mmol h^{-1} . The samples were mixed with the rare gas by sublimation between -10 and -15 °C in *high vacuo* (10^{-6} to 10^{-7} mbar). The concentrations of the samples in the matrices are unknown. Therefore, several experiments under various conditions for each isotopomer were conducted, to make sure that the molecules were indeed matrix isolated. A higher argon flow (8 mmol h^{-1}) and lower sublimation temperature (-20 °C) reduces the amount of the starting material significantly, but no differences in the IR spectra of the adducts and products were observed. Of course, at higher concentrations the IR spectra of the starting material change; i.e., the half band widths of the IR bands increase and new broad IR bands appear. That indicates that significant parts of the molecules are no longer matrix isolated under these conditions (e.g., sublimation temperatures of 0 °C).

The IR spectra of the matrices were recorded at 10 K on a Perkin-Elmer FTIR 1720x from 250 to 4000 cm^{-1} with a resolution of 1 cm^{-1} . A 450-W xenon arc lamp (Osram XBO 450W/4, PTI Inc. lamp housing A 5001) and a grating monochromator (PTI Inc. Model 001-01) were used for the irradiations. To obtain a high-energy output at 210 nm, the slit width of the monochromator was kept at 5 mm resulting in a band width of 20 nm.

The GAUSSIAN 92 package run on the Siemens-Nixdorf Vektorrechner S600/20 (Rechenzentrum der RWTH Aachen) was applied for *ab initio* calculations.⁷

All synthetic procedures were carried out under dry nitrogen in Schlenk glassware. Me_3Al (2.0 M in hexane) and $^{15}\text{NH}_3$ (98 atom % ^{15}N) were purchased from Aldrich and ND_3 (99.4 atom % D) from MSD Isotopes. NMR: Varian Unity 500 in C_6D_6 (24 °C) at 499.657 (^1H , standard: TMS intern), 130.195 (^{27}Al , standard: $[\text{Al}(\text{acac})_3]$ in C_6D_6 extern), and 125.639 MHz ($^{13}\text{C}\{^1\text{H}\}$, standard: TMS intern). MS: Varian Mat-CH-5 (70 eV). Elemental analyses (C, H, N): Heraeus Elementaranalysator CHN-O-RAPID.

$\text{Me}_3\text{Al}\cdot\text{NH}_3$. Gaseous NH_3 (250 mL at room temperature) was solidified with liquid nitrogen and 5.0 mL of AlMe_3 (2 M in hexane) was added. The cooling bath was removed, and the reaction mixture was stirred for 30 min at room temperature. After removal of the solvent in *vacuo*, sublimation (30 °C, 10^{-3} mbar) gave the pure product (0.86 g, 96%) as a colorless solid. ^1H NMR: δ -0.62 (s, 9H, CH_3), 0.26 (br s, 3H, NH_3). ^{13}C NMR: δ -8.06 (br). ^{27}Al NMR: δ 169 ($h_{1/2} = 1400$ Hz). MS: *m/e* 74 ($\text{M}^+ - 15, 100$). Anal. Calcd for $\text{C}_3\text{H}_{12}\text{AlN}$: C, 40.43; H, 13.57; N, 15.72. Found: C, 40.64; H, 13.93; N, 14.93.

$\text{Me}_3\text{Al}\cdot^{15}\text{NH}_3$. Following the procedure described for $\text{Me}_3\text{Al}\cdot\text{NH}_3$, $^{15}\text{NH}_3$ (500 mL) and 10.0 mL of AlMe_3 (2 M in hexane) yield 1.72 g of $\text{Me}_3\text{Al}\cdot^{15}\text{NH}_3$ (95%) after sublimation. ^1H NMR: δ -0.62 (s, 9H, CH_3), 0.26 (d, $J(^1\text{H}^{15}\text{N}) = 65.6$ Hz, 3H, NH_3). ^{13}C NMR: δ -8.06 (br). ^{27}Al NMR: δ 169 ($h_{1/2} = 1400$ Hz). MS: *m/e* 75 ($\text{M}^+ - 15, 100$). Anal. Calcd for $\text{C}_3\text{H}_{12}\text{Al}^{15}\text{N}$: C, 39.99; H, 13.42; N, 16.65. Found: C, 40.14; H, 13.43; N, 16.20.

$\text{Me}_3\text{Al}\cdot\text{ND}_3$. Following the procedure described for $\text{Me}_3\text{Al}\cdot\text{NH}_3$, ND_3 (250 mL) and 5.0 mL of AlMe_3 (2 M in hexane) yield 0.89 g of $\text{Me}_3\text{Al}\cdot\text{ND}_3$ (97%) after sublimation. ^1H NMR: δ -0.62 (s, CH_3). ^{13}C NMR: δ -8.15 (br). ^{27}Al NMR: δ 169 ($h_{1/2} = 1400$ Hz). MS: *m/e* 77 ($\text{M}^+ - 15, 100$). Anal. Calcd for $\text{C}_3\text{H}_9\text{D}_3\text{AlN}$: C, 39.11; N, 15.20. Found: C, 39.22; N, 14.56.

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